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# PVTx properties of the CO<sub>2</sub>-H<sub>2</sub>O and CO<sub>2</sub>-H<sub>2</sub>O-NaCl systems below 647 K: Assessment of experimental data and thermodynamic models

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#### Abstract

Evaluation of  $CO_2$  sequestration in formation brine or in seawater needs highly accurate experimental data or models of pressure–volume–temperature-composition (PVTx) properties for the  $CO_2$ – $H_2O$  and  $CO_2$ – $H_2O$ –NaCl systems. This paper presents a comprehensive review of the experimental PVTx properties and the thermodynamic models of these two systems. The following conclusions are drawn from the review: (1) About two-thirds of experimental data are consistent with each other, where the uncertainty in liquid volumes is within 0.5%, and that in gas volumes within 2%. However, this accuracy is not sufficient for assessing  $CO_2$  sequestration. Among the data sets for liquids, only a few are available for accurate modeling of  $CO_2$  sequestration. These data have an error of about 0.1% on average, roughly covering from 273 to 642 K and from 1 to 35 MPa; (2) There is a shortage of volumetric data of saturated vapor phase. (3) There are only a few data sets for the ternary liquids, and they are inconsistent with each other, where only a couple of data sets can be used to test a predictive density model for  $CO_2$  sequestration; (4) Although there are a few models with accuracy close to that of experiments, none of them is accurate enough for  $CO_2$  sequestration modeling, which normally needs an accuracy of density better than 0.1%. Some calculations are made available on www.geochem-model.org.

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Keywords: CO2 sequestration; PVTx properties; Volume; Density; Thermodynamic modeling

#### 1. Introduction

CO<sub>2</sub>-H<sub>2</sub>O and CO<sub>2</sub>-H<sub>2</sub>O-NaCl are the most commonly encountered fluids in various geological environments, such as the metamorphic rocks of almost all

grades, igneous rocks, carbonates, geothermal brines, sedimentary basins, and a wide variety of hydrothermal deposits (Roedder, 1984; Labotka, 1991; Schmidt and Bodnar, 2000). Therefore, the pressure–volume–temperature-composition (*PVTx*) and phase equilibrium properties of these two systems play a very important role in understanding petrologic and geochemical processes. Furthermore, the knowledge of the *PVTx* and

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phase equilibrium properties of these systems is also required by many processes in industrial production, and chemical and environmental engineering, such as CO<sub>2</sub>-related enhanced oil and gas recovery, supercritical fluid technology, seawater desalination and waste water treatment (Nighswander et al., 1989; Wang et al., 1996; Li et al., 2004; Ota et al., 2005; White et al., 2005). Today, with increasing interest in understanding the carbon cycle in Earth systems, and in storing CO2 in geological formations and sea waters, the PVTx and phase equilibrium properties of the systems CO<sub>2</sub>–H<sub>2</sub>O and CO<sub>2</sub>-H<sub>2</sub>O-NaCl are becoming increasingly important (Haugan and Drange, 1992; Oldenburg and Benson, 2002; Kaszuba et al., 2003; Baines and Worden, 2004; Pruess et al., 2004; Talman et al., 2004; Xu et al., 2004; Kervevan et al., 2005; Portier and Rochelle, 2005; Fuller et al., 2006).

It is already known that density differences about  $0.1~\rm kg~m^{-3}$  are sufficient to develop negatively buoyant plumes on kilometer scale in weakly stratified environments (Haugan and Drange, 1992), which is crucial for the natural transport of dissolved  $\rm CO_2$  to deeper water. Therefore, accurate prediction of liquid densities is very important. Under the P-T conditions of practical  $\rm CO_2$  sequestration, the total increase in solution density due to  $\rm CO_2$  dissolution are obviously larger than  $\rm 0.1~\rm kg~m^{-3}$ , but usually within a few percents of water density. On the other hand, the liquid density data sets in existence usually have an accuracy better than  $\rm 0.1\%$  (about 1.0 kg  $\rm m^{-3}$ ). Considering these facts, an accuracy better than  $\rm 0.1\%$  seems to be necessary for a density model of  $\rm CO_2$ -bearing solutions.

In this study, our interest focuses on the CO<sub>2</sub>–H<sub>2</sub>O and CO<sub>2</sub>–H<sub>2</sub>O–NaCl systems. We do not address the NaCl–H<sub>2</sub>O system, because its experimental *PVTx* and phase equilibrium properties and thermodynamic modeling have been studied extensively (Haas, 1970; Helgeson and Kirkham, 1976; Rogers and Pitzer, 1982; Pitzer and Pabalan, 1986; Bischoff et al., 1986; Li and Pitzer, 1987a,b; Pitzer et al., 1987; Pitzer and Tanger, 1988; Gallagher and Levelt Sengers, 1988; Tanger and Pitzer, 1989; Harvey and Levelt Sengers, 1989; Bischoff and Pitzer, 1989; Bischoff, 1991; Archer, 1992; Anderko and Pitzer, 1993).

The thermodynamic properties of  $CO_2$ – $H_2O$  and  $CO_2$ – $H_2O$ –NaCl have been studied for more than a century. However, there is still a lack of experimental data, even in some geologically and industrially important P–T regions, and the most recent models still cannot accurately represent PVTx and phase equilibrium properties in some P–T ranges. With the perspective of geochemical, environmental, and industrial applications

in mind, we have evaluated all available experimental *PVTx* data and the thermodynamic models for these two systems. Our goals were to identify the reliabilities of available data, to assess the strengths and weaknesses of the present thermodynamic models, and to identify important areas for future studies.

Comprehensive reviews on the CO<sub>2</sub> solubility or liquid-vapor equilibria have been published recently (Carroll et al., 1991; Crovetto, 1991; Duan et al., 1995; Diamond and Akinfiev, 2003; Duan et al., 2003; Duan and Sun, 2003; Spycher et al., 2003), and will not be repeated in this study. Tödheide (1982) reviewed the available PVTx data and some relevant thermodynamic models for the systems H<sub>2</sub>O, CO<sub>2</sub>-H<sub>2</sub>O, NaCl-H<sub>2</sub>O and CO<sub>2</sub>-H<sub>2</sub>O-NaCl, but at that time the data sets and thermodynamic models were limited. Ferry and Baumgartner (1987) summarized the equations of state for geological fluids before 1987, where the equations applicable for the binary include those of Holloway (1977), Kerrick and Jacobs (1981) and Spycher and Reed (1988). Mäder (1991) reviewed a few PVTx data sets for the CO<sub>2</sub>-H<sub>2</sub>O system relevant to mineral equilibria under high temperatures and pressures (Franck and Tödheide, 1959; Greenwood, 1969, 1973; Shmulovich et al., 1979, 1980; Gehrig, 1980; Sterner and Bodnar, 1991). In this study, we review the available experimental PVTx properties for both CO<sub>2</sub>-H<sub>2</sub>O and CO<sub>2</sub>-H<sub>2</sub>O-NaCl systems below the critical temperature of  $H_2O$ ,  $T_C(H_2O)$ .

# 2. Experimental PVTx properties of the $CO_2$ - $H_2O$ system

Published experimental PVTx data for the  $CO_2$ – $H_2O$  system are summarized in Table 1, and the P–T conditions for these measurements are plotted on Fig. 1.

# 2.1. The distribution of data in the P–T space and phase fields

Most data fall in the range of *P*<35 MPa (Ellis, 1959; Ellis and Golding, 1963; Parkinson and De Nevers, 1969; Zawisza and Malesińska, 1981; Wormald et al., 1986; Patel et al., 1987; Patel and Eubank, 1988; Nighswander et al., 1989; Crovetto and Wood, 1992; King et al., 1992; Ohsumi et al., 1992; Fenghour et al., 1996; Hnedkovsky et al., 1996; Teng et al., 1997; Zhang et al., 2002; Song et al., 2003b; Hebach et al., 2004). These data sets cover a temperature range from 276.15 to 644.78 K. Parkinson and De Nevers (1969) reported partial molar volumes of CO<sub>2</sub> in aqueous solutions, but did not give the corresponding pressures. The data set of

Table 1
The experimental *PVTx* properties of the CO<sub>2</sub>–H<sub>2</sub>O system

Authors	T/K	P/MPa	Concentration	$N_{ m d}$
Ohsumi et al. (1992)	276.15	34.754	$x(CO_2)\% = 0.1798 - 0.6294$	5
Song et al. (2003a,b)	273.25-284.15	5.0-12.5	$x(CO_2)\% = 0 - 0.6.1$	33
Teng et al. (1997)	278-293	6.44-29.49	$x(CO_2)\% = 2.50 - 3.49$	24
Parkinson and De Nevers (1969)	278.1-313.7	1.0342-34.4744	$x(CO_2)\% = 0.1 - 2.2$	28
Hebach et al. (2004)	283.80-333.19	1.09-30.66	$x(CO_2)$ = saturated values	201
King et al. (1992)	288.15-298.15	6.08-24.32	$x(CO_2)\% = 2.445 - 3.070$	27
Hnedkovsky et al. (1996)	298.15-633.41	1-35	$x(CO_2)\% = 0.279 - 0.332$	18
Zhang et al. (2002)	308.15	7.752-12.484	$x(CO_2)\% = 99.7$	16
Patel et al. (1987)	323.15-498.15	0.1-10	$x(CO_2) = 0.5 - 0.98$	423
Patel and Eubank (1988)	323.15-498.15	0.0855 - 10.0237	$x(CO_2) = 0.02 - 0.5$	297
Nighswander et al. (1989)	352.85-471.25	2.04-10.21	$x(CO_2)\% = 0.22 - 1.66$	33
Zawisza and Malesińska (1981)	373.15-473.15	0.385 - 3.35	$y(H_2O) = 0.1210 - 0.9347$	142
Ellis (1959)	387.15-621.15	0.5-16.4	$y(CO_2)\% = 3.90 - 6.27 - 84.02$	36
Fenghour et al. (1996)	415.36-644.78	5.884-27.964	$x(CO_2) = 0.0612 - 0.7913$	104
Wormald et al. (1986)	473.2-623.2	1-12	$x(CO_2) = 0.5$	57
Sterner and Bodnar (1991)	494.15-608.15	48.7-310	$x(CO_2) = 0.1234 - 0.7473$	84
Ellis and Golding (1963)	504.15-643.15	4.791-25.443	**	9
Singh et al. (2000)	573.15	7.44-99.93	$x(CO_2) = 0.05 - 0.9$	58
Seitz and Blencoe (1997)	573.15-623.15	99.93	$x(CO_2) = 0.1 - 0.9$	18
Blencoe et al. (2001b)	573.15	7.44-99.93	$x(CO_2) = 0 - 0.9$	148
Zakirov (1984)	573.15	5-180	$x(CO_2) = 0.3 - 0.77$	55
Zhang and Frantz (1992)	519.95-634.45	*	$x(CO_2)\% = 5.5 - 16.5$	29
Crovetto and Wood (1992)	622.75-642.70	19.64-28.13	$x(CO_2)\% = 0.48 - 0.8745$	72

Note:  $N_d$  = number of data points, x,y = mole fraction, \* pressures are not reported, but can be calculated from their empirical formulations, \*\* concentrations are not directly reported, but can be calculated from the report.

Zhang et al. (2002) is the only one for the dilute solution of H<sub>2</sub>O in supercritical CO<sub>2</sub> near the critical point of CO<sub>2</sub>. Song et al. (2003b) published graphical representation of their experimental densities and a simple empirical correlation of the ratio of solution and pure water densities. Hebach et al. (2004) gave the densities of aqueous solutions, but did not provide the corresponding composition information for the liquids, so their data should be used with a reliable solubility model. There are a few high-pressure data sets (up to 310 MPa) in the

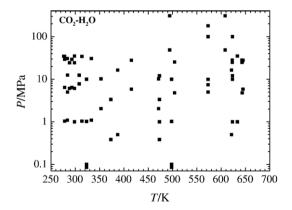


Fig. 1. The distribution of experimental PVTx data of the  $CO_2$ - $H_2O$  system in P-T space.

temperature range 494.15–608.15 K (Zakirov, 1984; Sterner and Bodnar, 1991; Seitz and Blencoe, 1997; Singh et al., 2000; Blencoe et al., 2001a).

A large portion of the data stated above were reported for a single saturated liquid or vapor phase (Parkinson and De Nevers, 1969; Zawisza and Malesińska, 1981; Nighswander et al., 1989; Sterner and Bodnar, 1991; King et al., 1992; Teng et al., 1997; Hebach et al., 2004). Few data sets are available for two coexistent phases (Ellis, 1959; Ellis and Golding, 1963).

Parkinson and De Nevers (1969) and Nighswander et al. (1989) presented CO<sub>2</sub>-saturated liquid density data. Zawisza and Malesińska (1981) measured dew-point densities at 373.15-473.15 K and 0.385-3.35 MPa. Sterner and Bodnar (1991) determined the molar volumes of synthetic fluid inclusions in vapor or liquid state at homogenization conditions (494.15-608.15 K and 47.7-310 MPa). King et al. (1992) measured the densities of both water-rich ( $L_{H,O}$ ) and  $CO_2$ -rich liquids  $(L_{CO_2})$  in equilibrium, but they did not find a significant difference between  $L_{\text{CO}_2}$  and pure  $\text{CO}_2$  liquid. Hebach et al. (2004) measured the densities of coexistent  $L_{\rm H,O}$ and  $L_{\text{CO}_2}$  phases in a wider P-T range. They confirmed the results of King et al. (1992) for the  $L_{\rm CO_2}$  phase, as can be expected from the very low content of water in it. Their measurements were conducted up to the point of

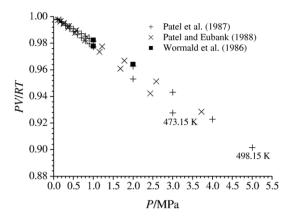


Fig. 2. Consistency between some gas PVTx data sets.

hydrate formation, but they did not report the equilibrium compositions. Teng et al. (1997) measured the densities of  $L_{\rm H,O}$  phases in equilibrium with  $L_{\rm CO2}$ . In some of their experiments ( $T \le 283$  K), CO<sub>2</sub> clathrate hydrate (H) formed at the  $L_{H,O}$ – $L_{CO}$ , interface, where the influence of hydrate on the measurements was considered to be negligible (Teng et al., 1997). However, as pointed out by Diamond and Akinfiev (2003) (p. 272), these measurements did not reach stable equilibrium. According to the phase rule, there are at most two phases co-existing away from the three-phase line, so  $H+L_{H_2O}+$  $L_{\rm CO_2}$  cannot stably coexist inside the hydrate region. We believe that the existence of hydrate in the experiments stopped the mass transfer between aqueous solution and liquid CO<sub>2</sub>. This is justified by the unusually sharp negative errors in CO<sub>2</sub> solubility found by Diamond and Akinfiev (2003) (p. 272).

There are numerous measurements of the formation conditions of  $CO_2$  hydrate, but almost none of them determined the corresponding densities and compositions. So far, no density data are available for the phase equilibria in the stable presence of hydrate.

## 2.2. The precisions, consistency, and reliability of data

The PVTx data of Seitz and Blencoe (1997), Singh et al. (2000) (a part) and Blencoe et al. (2001a) (a part) below  $T_{\rm C}({\rm H_2O})$  were obtained by the same technique as those of Seitz and Blencoe (1999). Seitz and Blencoe (1997) stated that their data are in excellent agreement with those of Fenghour et al. (1996) at 623.15 and 673.15 K. Additionally, Seitz and Blencoe (1999) have demonstrated that their data agree well with those of Shmulovich et al. (1980), Sterner and Bodnar (1991), and Fenghour et al. (1996). We found that the data of Seitz and Blencoe (1999) also agree well with the

molecular dynamics simulation results based on *ab initio* potentials Duan and Zhang (2006), which also agree well with other PVTx data at high temperatures and pressures. Considering these points, we believe that the above-mentioned data below  $T_{\rm C}({\rm H_2O})$  (Seitz and Blencoe, 1997; Singh et al., 2000; Blencoe et al., 2001a) are of good quality.

Fenghour et al. (1996) found that their data agree with those of Patel and Eubank (1988) within 1.5%, with an average deviation of about 0.4%. As mentioned just now, their data also agreement well with those of Seitz and Blencoe (1997).

Fig. 2 suggests that the gas phase *PVTx* data of Patel et al. (1987), Patel and Eubank (1988) and Wormald et al. (1986) agree very well with each other. This is remarkable considering the data of Wormald et al. (1986) were obtained from excess enthalpy measurements. On the other hand, if we extrapolate the data of Patel et al. (1987) to pure H<sub>2</sub>O, the results will be in good agreement with those of Wagner and Pruss (2002).

Fig. 3 shows that the data of Ohsumi et al. (1992), Hnedkovsky et al. (1996), Song et al. (2003b) and Li et al. (2004) are consistent with each other within an uncertainty of about 0.25%, where the average deviation is less than 0.1%. The water densities derived from the experimental data of Hnedkovsky et al. (1996) agree with those of Wagner and Pruss (2002) within 0.23% in most cases, with an exceptional point caused by calibration error (up to 2.12%) at 523.15 K and 19.99 MPa. The apparent molar volumes of Hnedkovsky et al. (1996) agree reasonably well with those of Crovetto and Wood (1992) if, as pointed out by Hnedkovsky et al. (1996), the sign in Crovetto's pressure correction is corrected. The pure water densities of Li et al. (2004)

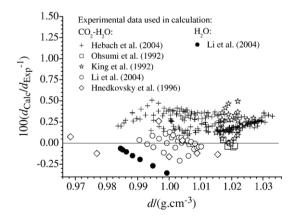


Fig. 3. Comparison of predicted and experimental densities of aqueous solutions of CO<sub>2</sub>. The solution densities are calculated with Eq. (1) (an empirical model of Song et al. (2003b) correlated from their own experimental data), and the water density is calculated with Eq. (2).

Table 2
The experimental *PVTx* properties of the CO<sub>2</sub>–H<sub>2</sub>O–NaCl system

Researchers	T/K	P/MPa	Concentration	$N_{ m d}$
Song et al. (2005)	276.15–283.15	4–13	wt.(CO <sub>2</sub> )%=0-7.7	64 <sup>A</sup>
Teng and Yamasaki (1998)	278–293	6.44-29.49	Sr%=3.5% $x(CO_2)\%=1.96-3.27$	$24^{\mathrm{B}}$
Song et al. (2003a,b)	303.15-323.15	10-20	m(NaCl) = 0.99 - 4.99 wt.(CO <sub>2</sub> )%=1-4	99 <sup>B</sup>
Li et al. (2004)	332.15	0.24-28.93	Sr=underground brine $c(CO_2)=0-0.958\times10^{-3}\text{mol}\cdot\text{cm}^{-3}$	37 <sup>C</sup>
Nighswander et al. (1989)	353.15-473.65	2.11-10.03	$Sr\%$ = Weyburn brine $x(CO_2)\% = 0.28-1.54$	34
Gehrig et al. (1986)	408-647	3-281.2	wt.(NaCl)%=1.0 $x(CO_2)$ %=0.18-84.9	52
Schmidt et al. (1995)	623.15	400	Sr% = 6.0-20.0 $x(CO_2)\% = 5$ (relative to H <sub>2</sub> O)	1
Krüger and Diamond (2001b)	540– <i>T</i> <sub>C</sub> (H <sub>2</sub> O)	65–350	Sr%=40 $x(CO_2)\%=9.69$ , $x(NaC1)\%=1.74$	Not given

A = a specific underground saline water within a Japanese city, B = artificial seawater, C = Weyburn brine, Sr = relative salinity, wt. = weight, m = molality, x = mole fraction, c = concentration.

have systematic deviations (up to 0.35%) from those of Wagner and Pruss (2002). The liquid densities of Hebach et al. (2004) agree with those of King et al. (1992) within 0.25%, with an average deviation of about 0.1%. Hebach et al. (2004) also reported that their data agree with those of Yaginuma et al. (2000) within 0.25% from 1 to 25 MPa, where the average deviation is less than 0.1%. However, if the data of Hebach et al. (2004) are extrapolated to the saturation pressures of pure water, the resulting densities will demonstrate systematic deviations (up to about 0.1–0.2%) from those of Wagner and Pruss (2002).

Ellis (1959) (p.233) stated that his results are not of high precision, where the vapor densities are unlikely to be more accurate than 5%. The sources of error include at least two factors: (1) the data do not truly represent the binary PVTx properties, because the CO<sub>2</sub>-H<sub>2</sub>O solutions contain both NaCl and NaHCO3; and (2) the suggested approach for the computation of solution volumes is approximate. The derived Henry's Law coefficients (K) of Ellis (1959) between 373.15 and 573.15 K are too low to be used for the CO<sub>2</sub> solubility calculation (Ellis and Golding, 1963). In this case, the CO<sub>2</sub> solubility calculation must use the distribution coefficients and vapor compositions. Ellis and Golding (1963) pointed out that their distribution coefficients derived from the liquid analysis are much better than those derived from the PVT method. If one wants to obtain more PVTx data, it is better to replace the distribution coefficients in their Table 5 with those fitted from their Table 2.

The density data of Parkinson and De Nevers (1969) seem to be questionable. The pressure dependence of

volume exhibited in their data at 278.1, 284.25, and 289.8 K are thermodynamically wrong. Their densities at 278.1 and 284.25 K decrease with increasing pressure (and  $CO_2$  concentration) after reaching a maximum, while the densities at 289.8 K remain constant after P>30 MPa. These trends are inconsistent with the other experimental results (King et al., 1992; Ohsumi et al., 1992; Hnedkovsky et al., 1996; Song et al., 2003b; Hebach et al., 2004; Li et al., 2004), in which the solution density increases with increasing pressure and  $CO_2$  concentration below 473.15 K. More importantly, the trends run against the rigorous theoretical

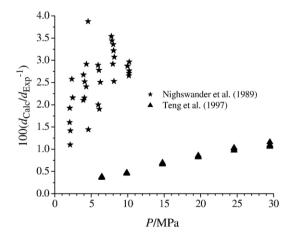


Fig. 4. Comparison of predicted and experimental densities of aqueous solutions saturated with CO<sub>2</sub>. The solution densities are calculated with Eq. (1) (an empirical model of Song et al. (2003b) correlated from their own experimental data), where the water densities are calculated with Eq. (2).

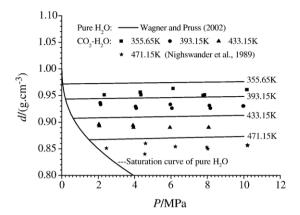


Fig. 5. Densities of saturated CO<sub>2</sub> solutions as a function of pressure at a given temperature. In principle, if the experimental densities at a given temperature are extrapolated to the saturated pressure of pure water at the same temperature, the resulting point should be on the saturation curve, but this is not the case for the data of Nighswander et al. (1989). That is, the experimental densities of Nighswander et al. (1989) have systematic deviations.

requirement of  $[\partial V/\partial P]_{T,x}$ <0 (which is also a common sense in thermodynamics).

The densities measured by Teng et al. (1997) are significantly lower than those of King et al. (1992) and Hebach et al. (2004) under the same conditions. They are also obviously inconsistent with other groups of lowtemperature density data. A comparison of Figs. 3 and 4 indicates that the densities determined by Nighswander et al. (1989) have large uncertainties and large negative deviations. This is evident if we extrapolate every group of isothermal density data to the saturated pressure of pure water (Fig. 5). In addition, the CO<sub>2</sub>-saturated densities determined by Nighswander et al. (1989) are almost constant as pressure (and CO<sub>2</sub> concentration) increases (Fig. 5). This is inconsistent with the accepted fact that the dissolution of CO<sub>2</sub> increases the liquid density, as demonstrated by the density measurements of Hnedkovsky et al. (1996) up to 473.35 K. Ji et al. (2005) also have similar findings. Therefore, the density data of Nighswander et al. (1989) and Teng et al. (1997) are unlikely to be of high precision.

Fig. 6 demonstrates that the gas *PVTx* data of Zawisza and Malesińska (1981) are inconsistent with those of Patel et al. (1987), and have larger deviations from predicted results of virial equation than those of Patel et al. (1987).

According to Zhang et al. (2002), their experimental pure  $CO_2$  volumes can be derived from their mixture densities and partial molar volumes of  $H_2O$ . The results deviate about -1% from those of Span and Wagner (1996) in the region far from the critical condition of

 $CO_2$ . The deviations become much larger (about -10%) near the critical conditions of  $CO_2$ .

Zakirov (1984) and Blencoe et al. (2001a) determined the excess volumes of  $CO_2$ – $H_2O$  at 573.15 K, where the results of Zakirov (1984) are systematically higher than those of Blencoe et al. (2001a) under the same conditions. The deviations decrease with increasing pressure from about 8 cm<sup>3</sup> mol<sup>-1</sup> to less than 1 cm<sup>3</sup> mol<sup>-1</sup>, which cause discrepancies of about 1–2% of the total volumes. As pressure increases, each isobaric excess volume-composition curve of Blencoe et al. (2001a) begins to show a negative minimum on the  $H_2O$ -rich side and a positive maximum on the  $CO_2$ -rich side. This variation is also followed by the spline fitting of Zakirov (1984) to his own experimental data, but his data at low  $x(CO_2)$ 's are scant (Fig. 7).

The solvus temperatures of CO<sub>2</sub>-rich phases determined by Sterner and Bodnar (1991) are 40 to 90 K higher than those of Takenouchi and Kennedy (1964), where the maximum disagreement occurs at about 150 MPa. For the same samples, Sterner and Bodnar (1991) reported a relatively large uncertainty in the total homogeneous temperatures (T<sub>h</sub>), which is about 6 K for total homogenization to liquid, and up to about 25 K for total homogenization to vapor. Blencoe et al. (2001b) reported new vapor-liquid equilibrium data for the binary at 573.15 K, which agree well with those of Takenouchi and Kennedy (1964), and are claimed to be highly accurate. These data can be used to examine the reliability of the solvus data of Sterner and Bodnar (1991). According to these new data and other previous solvus data, we think that the errors in the solvus data of Sterner and Bodnar (1991) are quite large. This is not due to the underestimated  $T_h$ 's of the fluid inclusions that homogenized to vapor, but rather to the solvus pressures calculated from the equation of

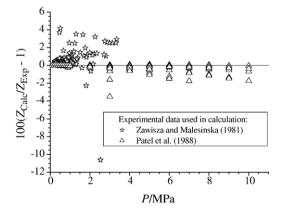


Fig. 6. Deviations of predicted compressibility factors of gaseous CO<sub>2</sub>–H<sub>2</sub>O mixtures. Compressibility factors are calculated with the cubic virial equation of state of Patel et al. (1987).

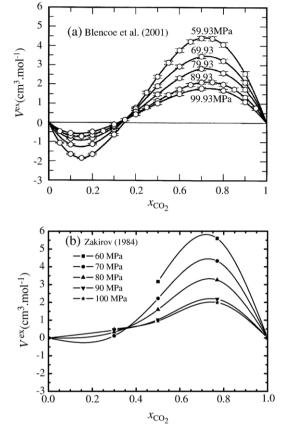


Fig. 7. Excess volumes of CO<sub>2</sub>-H<sub>2</sub>O mixtures at 573.15 K.

Connolly and Bodnar (1983). The solvus pressures were overestimated up to about 30–70 MPa for the 543.15 K isotherm, and up to about 5–11 MPa for the 573.15 K isotherm (Blencoe et al., 2001b). Therefore, the solvus pressures of Sterner and Bodnar (1991) should be corrected with a reliable equation of state or other approaches.

The *P*–*T* curves of the binary vapor-liquid equilibria obtained by Zhang and Frantz (1992) are significantly different (about 20–70 K or 10–50 MPa) from those of Takenouchi and Kennedy (1964). Similar differences also exist between the results of Zhang and Frantz (1992) and those of Tödheide and Franck (1963), Seitz and Blencoe (1997) and Blencoe et al. (2001a,b). The high accuracy of the data of Blencoe and co-workers suggests that the vapor-liquid equilibrium data of Zhang and Frantz (1992) may have systematic deviations.

# 3. Experimental *PVTx* properties of the CO<sub>2</sub>–H<sub>2</sub>O–NaCl system

As shown in Table 2, much less experimental *PVTx* data are available for this system than those for the

CO<sub>2</sub>-H<sub>2</sub>O system. The data distribution, and precision, and the agreement among data sets are discussed below.

## 3.1. Distribution of data in P-T space and phase fields

Several experimental data sets are reported for this system (Gehrig, 1980; Gehrig et al., 1986; Nighswander et al., 1989; Schmidt et al., 1995; Teng and Yamasaki, 1998; Krüger and Diamond, 2001b; Song et al., 2003a; Li et al., 2004; Song et al., 2005). Some high-pressure data were reported by Gehrig (1980), Gehrig et al. (1986), Krüger and Diamond (2001b), and Schmidt et al. (1995), where the data of Gehrig et al. (1986) are revised from those of Gehrig (1980). Schmidt et al. (1995) have only one datum point below  $T_{\rm C}({\rm H_2O})$ . The reported salinities of Gehrig (1980), Gehrig et al. (1986), and Schmidt et al. (1995) are up to 20 and 40 wt.% (relative to H<sub>2</sub>O), respectively, and the CO<sub>2</sub> contents are up to 84.9 mol.%. The other data sets are limited to low pressure (<30 MPa), low CO<sub>2</sub> content (<10 mol.%) and low salinity (<8 mol.%). Most of these data are in the range 273.15-333.15 K, where the measurements of Nighswander et al. (1989) from 355.15 K to 473.65 K are exceptional. The data distribution is plotted in Fig. 8. Apparently, there are no data for T<400 K and P>30 MPa.

The *PVTx* data for the homogeneous mixtures include those of Schmidt et al. (1995) and Song et al. (2003a, 2005). Song et al. (2003a) measured the densities of CO<sub>2</sub>-bearing underground brine, but they did not mention the corresponding salt contents. Song et al. (2005) reported density diagrams of artificial seawater at 4–13 MPa, 276.15 and 283.15 K.

The data sets on phase equilibria include those of Gehrig (1980), Gehrig et al. (1986), Krüger and Diamond (2001b), Nighswander et al. (1989), Teng and

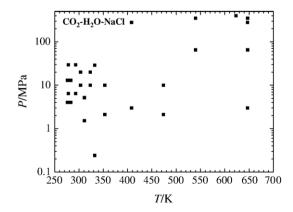


Fig. 8. Distribution of experimental PVTx data of the  $CO_2$ - $H_2O$ -NaCl system in P-T space.

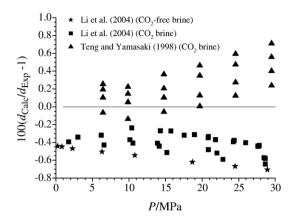


Fig. 9. Deviations of predicted densities of brines containing  $CO_2$ . The mixture densities are calculated with the model of Song et al. (2005), whereas the densities of the reference system NaCl-H<sub>2</sub>O are calculated with the model of Rogers and Pitzer (1982), and the water densities are calculated with our Eq. (2).

Yamasaki (1998), and Li et al. (2004), all reporting saturated liquid volumes or densities. Gehrig and coworkers reported saturated vapor volumes for  $T \le T_{\rm C}$  (H<sub>2</sub>O) (Gehrig, 1980; Gehrig et al., 1986). It should be noted that at temperatures below 573K the solubility of NaCl in the vapor phase is very low, so it is a good approximation in practice to treat the vapor phase as a  ${\rm CO_2-H_2O}$  mixture.

Gehrig and coworkers presented extensive data on volumes of saturated liquid and vapor phases, covering a range of 408–793 K and 3–281.2 MPa (Gehrig, 1980; Gehrig et al., 1986). Krüger and Diamond (2001b) presented a graphical report of their experimental volumes on the bubble and dew curves and the relevant iso- $T_h$  lines (contours of total homogenization temperatures). According to experimental synthetic fluid inclusion data, Schmidt et al. (1995) constructed a P-T projection of the dew curve of the mixture containing 40 wt.% NaCl (relative to  $H_2O$ ) and 5 wt.%  $CO_2$ . It partially falls within the range of  $T \le T_C(H_2O)$ .

Nighswander et al. (1989) and Li et al. (2004) measured the densities of saline water saturated with gaseous CO<sub>2</sub>, where the density data of Li et al. (2004) were determined for the Weyburn Formation brine. The experiments of Teng and Yamasaki (1998) were conducted for synthetic seawater in the presence of liquid CO<sub>2</sub> at 278–293 K, where some of their experimental conditions fall within the CO<sub>2</sub> hydrate formation region, for which the influence of hydrate was considered to be negligible. Ohmura and Mori (1999) analyzed the temperature dependence of solubility measured by Teng and Yamasaki (1998), Stewart and Munjal (1970), and Kimuro et al. (1994), and concluded that the data of

Teng and Yamasaki (1998) represent the approximate *PVTx* properties of the system in the metastable absence of hydrate, not those of the seawater in equilibrium with hydrate. As concluded by Diamond and Akinfiev (2003) for the data measured by Teng et al. (1997), the corresponding experiments did not reach stable equilibrium states. We believe that the existence of hydrate in the experiments of Teng and Yamasaki (1998) stopped the mass transfer between aqueous solution and liquid CO<sub>2</sub>. It is important to point out that no volumetric data are available for the fluids in stable equilibrium with CO<sub>2</sub> hydrate, although there have been numerous measurements of CO<sub>2</sub> hydrate formation conditions for the ternary system.

#### 3.2. Precision, consistency, and reliability of data

Few comments on the precision or consistency can be found for the ternary *PVTx* data. Song et al. (2003a) noted that the measurements in literature, e.g., Nighswander et al. (1989), Ohsumi et al. (1992), and their previous determinations, have respective shortcomings in their measurement principles and precisions, but they failed to provide further comments.

Like the binary density data discussed earlier, the ternary density data of Nighswander et al. (1989) also have large negative deviations and large uncertainties. Fig. 9 suggests that the density data of Teng and Yamasaki (1998) and Li et al. (2004) are inconsistent with those of Song et al. (2005). It should be noted that the densities of NaCl–H<sub>2</sub>O solutions of Li et al. (2004) deviate by about 0.4%–0.7% from the results of Rogers and Pitzer (1982) model. These deviations are greater than the error range of the model. In addition, Wong et al. (2005) found that the salting-out coefficients from

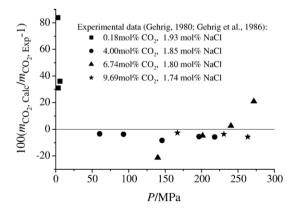


Fig. 10. Comparison of predicted CO<sub>2</sub> solubilities with experimental results. The solubilities (in molality) are calculated with the model of Duan et al. (2006).

the equation suggested by Teng and Yamasaki (1998) have two distinct features: (1) The coefficients decrease with increasing temperature; (2) They become positive close to 273 K. The first feature contradicts the salting-out coefficients of Weiss (1974), which were considered to be well-established (Wong et al., 2005). The second feature contradicts all the experimental data for CO<sub>2</sub> solubility in seawater (Wong et al., 2005).

In order to resolve the inconsistency problem in some PVTx data in literature, Krüger and Diamond (2001b) examined the immiscibility boundaries and the relevant iso-T<sub>h</sub> lines of two CO<sub>2</sub>-H<sub>2</sub>O-NaCl mixtures (88.57 mol.% H<sub>2</sub>O+9.69 mol.% CO<sub>2</sub>+1.74 mol.% NaCl; 77.47 mol.% H<sub>2</sub>O+21.01 mol.% CO<sub>2</sub> and 1.52 mol.% NaCl) using the synthetic fluid inclusion technique. Below T<sub>C</sub>(H<sub>2</sub>O), their bubble curve of the first mixture has higher pressures than that of Schmidt and Bodnar (2000), but it coincides perfectly with those of Gehrig (1980), and is very close to the points of Takenouchi and Kennedy (1965). The bubble-point curve of the first mixture reported by Gehrig (1980) is also supported by the solubility model of Duan and Sun (2003) (Fig. 10). For the second mixture, however, the dew-point curve of Krüger and Diamond (2001b) is different from that of Gehrig (1980). The bubble-point compositions of the second mixture determined by Gehrig (1980) below  $T_C(H_2O)$  also have large deviations from the predicted results of solubility model Duan and Sun (2003) (Fig. 10). The largest deviations are notably out of the error range of the model. These results suggest that some of the bubble-point and dewpoint curves determined by Gehrig (1980) are unreasonable. As for bubble-point volumes, the results of the first mixture obtained by Gehrig (1980) are notably

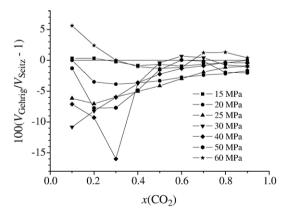


Fig. 11. Difference between the experimental volumes of  $\rm CO_2$ – $\rm H_2O$  mixtures at 673.15 K determined by Gehrig (1980) and Seitz and Blencoe (1999).  $V_{\rm Gehrig}$  and  $V_{\rm Seitz}$  are the volumes determined by Gehrig (1980) and Seitz and Blencoe (1999), respectively.

lower than those of Krüger and Diamond (2001b), where the maximum deviation is up to about 10% around the critical point. Much greater discrepancy is found between the dew-point volumes of the second mixture. In addition, the volumes of CO<sub>2</sub>–H<sub>2</sub>O obtained by Gehrig (1980) are also notably lower than those of Seitz and Blencoe (1999) (Fig. 11). Similarly, the volumes of NaCl–H<sub>2</sub>O of Gehrig (1980) are also lower than the predicted results of Rogers and Pitzer (1982) model, where the differences is obviously out of the possible error range of the model. In summary, the volume data of Gehrig (1980) have large negative deviations.

# 4. Thermodynamic modeling of the mixture PVTx properties

As reviewed in the last section, the experimental PVTx data of the  $CO_2$ – $H_2O$  and  $CO_2$ – $H_2O$ –NaCl systems only occupy a very limited P–T range as compared to the vast P–T range encountered in the Earth's interior. Therefore, the modeling of PVTx and phase equilibrium properties of the two systems that allows accurate extrapolation over a wide P–T range is extremely important. Many thermodynamic models have been developed to predict volumetric properties of the two systems. They can be divided into five types according to the features of volumetric properties: (1) equations of state (EOS); (2) liquid density models; (3) partial molar volume models; (4) apparent molar volume models; and (5) excess volume models. All these models are assessed in this report.

## 4.1. The $CO_2$ – $H_2O$ system

## 4.1.1. Equations of state

At low to medium pressures, numerous cubic EOS and virial EOS truncated at the second or third virial coefficient are adequate to predict the PVTx, phase equilibria and other thermodynamic properties of fluids (including CO<sub>2</sub>–H<sub>2</sub>O). Thus, these two types of EOS are widely used in physical chemistry, chemical engineering and geochemistry (Shmulovich et al., 1980; Vanderzee and Haas, 1981; Patel et al., 1987; Wormald and Lloyd, 1994; Plyasunov and Shock, 2003). For the CO<sub>2</sub>–H<sub>2</sub>O system, there are many correlations of the second and third virial coefficients in literature (Tsonopouplos, 1974; Holste et al., 1987; Hendl et al., 1997; Abdulagatov et al., 2002; Plyasunov and Shock, 2003; Harvey and Lemmon, 2004). However, these equations cannot be directly extended to the high-pressure fluids that are frequently encountered in many natural systems. For example, Wormald and Lloyd (1994) used a cubic EOS for excess enthalpies of CO<sub>2</sub>—H<sub>2</sub>O gas mixtures up to 913 K and 20 MPa. The predicted enthalpies agree well with experimental results up to 10 MPa, but the deviations from experimental results increase rapidly with pressure. For the CO<sub>2</sub>—H<sub>2</sub>O system, cubic virial equation can be applied to about 10 MPa (Fig. 6), and the quadratic equation is only valid up to a few MPa. For these reasons, many EOS have been developed for the CO<sub>2</sub>—H<sub>2</sub>O mixtures in a wider *P*–*T* range (Kerrick and Jacobs, 1981; Connolly and Bodnar, 1983; Nitsche et al., 1984; Spycher and Reed, 1988; Duan et al., 1992b; Churakov and Gottschalk, 2003b; Ji et al., 2005).

Nitsche et al. (1984) used a lattice-gas model to represent the behavior of CO<sub>2</sub>, H<sub>2</sub>O, and their mixtures. This model correctly predicts the qualitative patterns of mixture phase behavior using two mixture parameters fitted to simple temperature functions over a small temperature range. They claimed that quantitative prediction of mixture phase behavior is possible over a not too wide temperature range.

The virial equation of Spycher and Reed (1988) is truncated at the third virial coefficient, with *P* and *T* as variables. It was claimed to be valid from 323.15 to 623.15 K, and from 0 to 50 MPa. However, this equation underestimates the volumes of gaseous mixtures. The deviations of predicted volumes from experimental results (Wormald et al., 1986; Patel et al., 1987; Fenghour et al., 1996) increase with increasing pressure, where the average deviation is about 4%, and the maximum deviation is nearly 10%.

The equation of Duan et al. (1992b) can calculate both phase equilibria and volumetric properties of the CO2-H2O system with accuracies close to those of experiments at 323-1273 K and 0-100 MPa (or 300 MPa with less accuracy). For the end-members, it is valid at 273-1273 K and 0-800 MPa (or slightly above), but it cannot accurately predict the phase equilibria of end-members near the triple points. For the binary, the equation cannot correctly reproduce nearcritical PVTx properties and phase behavior (Seitz and Blencoe, 1999; Bakker and Diamond, 2000; Blencoe, 2004). Because of the high overall accuracy, the equations of Duan et al. (1992a,b) were incorporated into a software package HCO-TERNARY written by Nieva and Barragán (2003). However, the accuracy of this equation cannot be better than 0.25% in volume, so it is still inadequate for the modeling of CO<sub>2</sub> sequestration.

Churakov and Gottschalk (2003a,b) proposed a perturbation theory-based EOS for both polar and non-polar molecular fluids and their mixtures. The CO<sub>2</sub>–H<sub>2</sub>O phase boundaries predicted with this equation deviate significantly from experimental results (Tödheide and

Franck, 1963; Takenouchi and Kennedy, 1964), especially near and in the critical regions. Generally, the equation can reasonably reproduce experimental excess volumes (Gehrig, 1980; Shmulovich et al., 1980; Sterner and Bodnar, 1991; Frost and Wood, 1997; Seitz and Blencoe, 1999), where the maximum deviations range from less than 2 cm<sup>3</sup>·mol<sup>-1</sup> to more than 10 cm<sup>3</sup>·mol<sup>-1</sup>. Churakov and Gottschalk (2003a,b) did not report the performance of the EOS in binary PVTx prediction below  $T_{\rm C}({\rm H_2O})$ , which should not be better than the performance above  $T_{\rm C}({\rm H_2O})$  because of the greater difficulty in PVTx modeling below  $T_C(H_2O)$ . Considering the deviations in excess volume from experimental data, we believe that the error of this equation should be more than 0.5%, so it is not applicable for CO<sub>2</sub> sequestration.

Ji et al. (2005) used a semi-empirical statistical mechanics model (SAFT1-RPM) to model the densities and phase equilibria of the CO<sub>2</sub>-H<sub>2</sub>O system in the range 285-473 K, 0-60 MPa and the CO<sub>2</sub>-H<sub>2</sub>O-NaCl system in the range 298-373 K and 0-20 MPa. The model can reasonably reproduce both vapor and liquid density data of the binary (Patel and Eubank, 1988; King et al., 1992; Hebach et al., 2004; Li et al., 2004), where the average deviations of predicted liquid densities are about 0.35%. This model can also predict the solubilities of CO<sub>2</sub> in water and aqueous NaCl solutions, but its accuracy is generally inferior to the model of Duan et al. (2006). Because of the lack of reliable ternary density data, the accuracy of the model in ternary density prediction remains to be examined, even though it is not expected to be better than that of the binary.

## 4.1.2. Liquid density models

Some density models of aqueous CO<sub>2</sub> solutions can be regarded as another type of EOS of CO<sub>2</sub>–H<sub>2</sub>O (Teng et al., 1997; Garcia, 2001; Bachu and Adams, 2003; Song et al., 2003b). Teng et al. (1997) presented an approximate density difference model to correlate the densities of aqueous CO<sub>2</sub> solutions at low temperatures and low to medium pressures, where the parameters were taken as P-T independent. This kind of model is only valid in a small P-T range. The density model of Garcia (2001) was based on an approximate apparent molar volume expression regressed from experimental data in the range 278.15-573.15 K. It neglects the effects of pressure and CO<sub>2</sub> concentration on apparent molar volume. At low temperatures, the apparent molar volumes predicted with this model are inconsistent with other data (Ohsumi et al., 1992; Song et al., 2003b) due to the use of inaccurate density data of Teng et al. (1997) in the parameter regression. The CO<sub>2</sub>-H<sub>2</sub>O density

model of Bachu and Adams (2003) consists of the water density model of Batzle and Wang (1992) and the apparent molar volume expression of Garcia (2001), where the water model deviates from the well-established results of Wagner and Pruss (2002) up to 0.2% below 373.15 K, and increases rapidly with increasing temperature and decreasing pressure. Song et al. (2003b) proposed a density model for the dilute solution of CO<sub>2</sub>:

$$\rho_{\rm aq} = \rho_w [1 + 0.275 wt({\rm CO}_2)] \tag{1}$$

where  $\rho_{\rm aq}$ ,  $\rho_{\rm w}$  and  $wt({\rm CO_2})$  are solution density, water density and mass fraction of  ${\rm CO_2}$ , respectively. However, they did not give the corresponding  $\rho_{\rm w}$  expression. In order to facilitate its application, we develop a formula for the molar volumes of liquid water:

$$V_{w} = k_{0} + k_{1}P + k_{2}P^{2} + k_{3}P^{3}$$

$$k_{i} = k_{i1}T^{3} + k_{i2}T^{2} + k_{i3}T + k_{i4} + k_{i5}T^{-1} \quad (i = 0, 1)$$

$$k_{j} = k_{j1}T^{3} + k_{j2}T^{2} + k_{j3} \quad (j = 2, 3)$$
(2)

where the units of  $V_W$ , T and P are cm<sup>3</sup>•mol<sup>-1</sup>, K and bar, respectively, and the constants are given in Table 3. Eq. (2) reproduces the water volumes of Wagner and Pruss (2002) with an average deviation of 0.005% and a maximum deviation of 0.027% in the range 273.15–473.15 K and 0–200 MPa (Fig. 12).

#### 4.1.3. Partial molar volume models

Current thermodynamic description of the aqueous non-electrolyte solutions is far less accurate than those for electrolyte solutions (Akinfiev and Diamond, 2003, 2004). For example, the revised equation of Helgeson, Kirkwood and Flowers (HKF) (Tanger and Helgeson, 1988; Shock et al., 1989; Shock and Helgeson, 1990; Shock et al., 1992) is designed to predict the thermodynamic properties of aqueous ions over a wide *P*–*T* range (273–1273 K and 0.1–500 MPa) (Akinfiev and Diamond, 2003). However, some deficiencies are found when it is extended to aqueous non-electrolytes. For example, it does not correctly predict the behavior of

Table 3 Constants in Eq. (2)

$k_{01}$	3.27225E-07	$k_{14}$	3.59860E-02
$k_{02}$	-4.20950E-04	$k_{15}$	-3.55071E+00
$k_{03}$	2.32594E-01	$k_{21}$	2.57241E-14
$k_{04}$	-4.16920E+01	$k_{22}$	-1.24336E-11
$k_{05}$	5.71292E+03	$k_{23}$	5.42707E-07
$k_{11}$	-2.32306E-10	$k_{31}$	-4.42028E-18
$k_{12}$	2.91138E-07	$k_{32}$	2.10007E-15
$k_{13}$	-1.49662E-04	$k_{33}$	-8.11491E-11

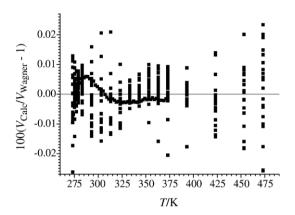


Fig. 12. Deviations of water volumes predicted with Eq. (2).  $V_{\rm Wagner}$  is the molar volume predicted with the model of Wagner and Pruss (2002).

non-electrolyte solutes in the near-critical and supercritical regions of water (O'Connell et al., 1996; Plyasunov et al., 2000a; Plyasunov and Shock, 2001; Akinfiev and Diamond, 2003). Akinfiev and Zotov (1999) demonstrated that the HKF model cannot correctly predict the activity coefficients of aqueous  $CO_2$  at infinite dilution below 573 K. This indicates that the corresponding excess partial volumes are also inaccurate. The infinite dilution partial molar volumes of  $CO_2$  ( $\overline{V}_{CO_2}^{\infty}$ ) predicted with the HKF model (Shock et al., 1989) have large deviations from the experimental results of dilute solutions of  $CO_2$  (Crovetto and Wood, 1992; Hnedkovsky et al., 1996). Thus, we need to develop equations for the partial molar volumes of aqueous nonelectrolytes.

Akinfiev and Diamond (2003, 2004) proposed a three-parameter equation for the volumetric properties of aqueous non-electrolytes at infinite dilution, where the predicted  $(V_{CO_2}^{\infty})$ 's of  $CO_2$  are in reasonable agreement with experimental results. Plyasunov et al. (2000a,b) and Sedlbauer et al. (2000) briefly reviewed the progress in the modeling of volumetric properties of aqueous non-electrolyte solutions, where one of the most important advances is the approach based on the fluctuation solution theory (Kirkwood and Buff, 1951; O'Connell, 1971; O'Connell, 1981; O'Connell, 1995; O'Connell et al., 1996; Harvey, 1998; Sedlbauer et al., 2000; Plyasunov et al., 2000a,b, 2006), which can be used to develop accurate  $\overline{V}_{\text{CO}}^{\infty}$ , models over a wide P-Trange, including near-critical and supercritical regions (Harvey, 1998; Plyasunov et al., 2000b). However, if these  $V_{\text{CO}}^{\infty}$ , models are used to predict the densities or total volumes of CO2 solutions, they should be confined to the dilute solutions of CO<sub>2</sub>, because the dependence of volume on composition is highly non-linear.

#### 4.1.4. Apparent molar volume models

In principle, apparent molar volume  $(V_{\phi})$  models of aqueous  $CO_2$  can be directly used to predict the total volumes  $(V_{aq})$  or densities of  $CO_2$ – $H_2O$  solutions:

$$V_{\text{aq}} = x_{\text{H}_2\text{O}} V_{\text{H}_2\text{O}} + x_{\text{CO}_2} V_{\phi,\text{CO}_2} \rho_{\text{aq}} = (x_{\text{H}_2\text{O}} M_{\text{H}_2\text{O}} + x_{\text{CO}_2} M_{\text{CO}_2}) / (x_{\text{H}_2\text{O}} V_{\text{H}_2\text{O}} + x_{\text{CO}_2} V_{\phi,\text{CO}_2})$$
(3)

where M is relative molecular weight, x is mole fraction, and V is molar volume. If these  $V_{\phi}$  models are combined with appropriate NaCl-H<sub>2</sub>O volume (or density) models and ternary excess volume models, the total volumes (or densities) of CO<sub>2</sub>-H<sub>2</sub>O-NaCl solutions can also be predicted.

Presently, there are some empirical or semi-theoretical models for  $V_{\phi, \text{CO}_2}$  (Crovetto et al., 1990; Andersen et al., 1992; Enick and Klara, 1992; Iglesias and Moya, 1992; Garcia, 2001). Except for the model of Crovetto et al. (1990), these models neglected the effects of pressure and composition on apparent molar volumes, so they cannot accurately reproduce experimental results (Fig. 13). The model of Crovetto et al. (1990) was proposed to correlate  $V_{\phi, \text{CO}_2}$  in near-critical region of H<sub>2</sub>O. Therefore, they employed the current near-critical theories of dilute solution that are believed to be successful in the treatment of critical singularity. This method is also applicable below  $T_{\text{C}}(\text{H}_2\text{O})$ , but no ana-

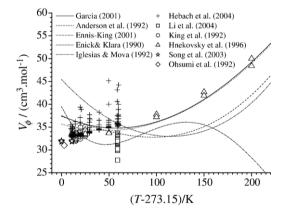


Fig. 13. Predicted apparent molar volumes of  $CO_2$  in water compared with experimental results. The apparent molar volume model of Ennis-King (2001) was given by Garcia (2001) through his personal communication with Jonathan Ennis-King. The apparent molar volumes of Hnedkovsky et al. (1996) are directly given by the original authors. The other apparent molar volumes are calculated with experimental densities, where the water densities are calculated with Eq. (2). The  $CO_2$  solubilities corresponding to the saturated density data of Li et al. (2004) and Hebach et al. (2004) are calculated with the model of Duan et al. (2006).

lytical expression was given for its parameters, except for single parameter values under given P-T conditions.

#### 4.1.5. Excess volume models

Some excess volume models have been proposed for the  $CO_2$ – $H_2O$  system, but most of them are developed for elevated pressures and temperatures (Shmulovich et al., 1980; Sterner and Bodnar, 1991; Destrigneville et al., 1996; Blencoe et al., 1999). Few models are developed for excess volumes of the binary below  $T_C(H_2O)$ .

In addition to the direct prediction of excess volumes, there are also many other ways for the indirect derivation of volumes. For example, excess volumes can be derived from equations of excess enthalpy, or excess free energy. These approaches have been used for the CO<sub>2</sub>-H<sub>2</sub>O system, such as the two-fluid free energy model (Gallagher et al., 1993; Wormald et al., 1997) and the excess enthalpy models (Wormald et al., 1983; Smith and Wormald, 1984; Bottini and Saville, 1985; Wormald et al., 1986). In fact, only a small part of the numerous excess function models in physical chemistry and chemical engineering have been introduced into the thermodynamic modeling of geological fluids, and most of them are old, such as Margules-type, Van Laar-type, Wilson-type (Wilson, 1964) and Non-Random Two-Liquid (Renon and Prausnitz, 1968) models. Better excess function models for geological fluids can be developed, if we use more advanced thermodynamic theories. Many theoretical or semi-empirical local composition and coordination number models can serve as a good base, such as those of Aranovich and Donohue (1996), Aranovich et al. (1997), Cao and Wang (2000), Guo et al. (1990), and Hu and Duan (2005).

#### 4.2. The CO<sub>2</sub>-H<sub>2</sub>O-NaCl system

A few models on the PVTx and phase equilibria have been proposed for this system. Some of them are developed for temperatures above 573 K or 623 K (Bowers and Helgeson, 1983, 1985; Brown and Lamb, 1989; Duan et al., 1995), and are usually inadequate below  $T_{\rm C}({\rm H_2O})$ . The other models are developed to meet the needs of enhanced oil and gas recovery and  $CO_2$  sequestration at low to medium P-T conditions (Teng and Yamasaki, 1998; Bachu and Adams, 2003; Bando et al., 2004; Ji et al., 2005; Song et al., 2005), and most of them are simple approximate density models. For example, Teng and Yamasaki (1998) and Song et al. (2005) used a density difference expression to correlate the densities of CO<sub>2</sub>-bearing seawater. Bando et al. (2004) used the density difference expression of Teng et al. (1997) for CO<sub>2</sub>-H<sub>2</sub>O solutions and the density

model of Pitzer et al. (1984) for NaCl-H<sub>2</sub>O solutions to predict the densities of CO<sub>2</sub>-NaCl-H<sub>2</sub>O solutions. The three density difference expressions are all regressed from densities of dilute CO<sub>2</sub> solutions, and ignore the effects of pressure and temperature.

Different from these approaches, Bachu and Adams (2003) proposed an approximate model to predict the CO<sub>2</sub>-bearing brine densities, where the apparent molar volumes of CO2 in water is calculated from the approximate expression of Garcia (2001), the densities of water and NaCl solutions are calculated with the empirical models of Batzle and Wang (1992). We found that the NaCl solution densities predicted from the model of Batzle and Wang (1992) are systematically (0.45% on average) lower than those from the model of Rogers and Pitzer (1982). Because of the use of multiple approximations, the densities predicted with this model have significant negative deviations (-0.87% on average) from those predicted with the model of Song et al. (2005), where the latter model can accurately reproduce the experimental results of Song et al. (2005).

#### 5. Conclusions

## 5.1. Experimental PVTx data for $CO_2$ – $H_2O$ and $CO_2$ – $H_2O$ –NaCl fluids

For  $\mathrm{CO_2-H_2O}$  mixtures, the experimental PVTx data range from 276.15 to 644.78 K, and from 0.0855 to 310 MPa. New data are needed to fill the experimentally blank region of T<493 K and P>35 MPa. Below  $T_{\mathrm{C}}$  (H<sub>2</sub>O), a considerable portion of the binary volumetric data are measured for single saturated phases (liquid or vapor), and most of them are for saturated liquids. There is a shortage of saturated vapor volumes and the  $\mathrm{CO_2-rich}$  liquid volumes at low temperatures. In the region close to  $T_{\mathrm{C}}(\mathrm{H_2O})$  and  $P_{\mathrm{C}}(\mathrm{H_2O})$ , accurate or reliable PVTx data are scarce, especially for the compositions with high  $\mathrm{CO_2}$  contents.

About one-third of the reported data sets (Ellis, 1959; Ellis and Golding, 1963; Zawisza and Malesińska, 1981; Nighswander et al., 1989; Sterner and Bodnar, 1991; Zhang and Frantz, 1992; Teng et al., 1997; Zhang et al., 2002) have large systematic errors or uncertainties, or are significantly inconsistent with other data sets. Inconsistencies between these data and the others below  $T_{\rm C}({\rm H_2O})$  can be very large, especially those between the PVTx data along phase boundaries. If we exclude the questionable data sets, the differences between the vapor volume data below  $T_{\rm C}({\rm H_2O})$  are generally within 2%, and the liquid density data are within 0.5% for most cases. The phase boundary data usually have larger

deviations. The discrepancies between data sets increase considerably with increasing temperature and pressure, indicating increasing uncertainties in *PVTx* measurements at higher temperatures and pressures.

For the  $CO_2$ – $H_2O$ –NaCl system, the experimental PVTx data cover a temperature range from 276.15 to 647 K and a pressure range from 0.24 to 400 MPa. No PVTx data are available for T<400 K and P>30 MPa. Almost no accurate density measurements are made for the phase equilibria with  $CO_2$  hydrate, although there are numerous studies of  $CO_2$ – $H_2O$ –NaCl fluid inclusions at low temperatures.

Several *PVTx* data sets for CO<sub>2</sub>–H<sub>2</sub>O–NaCl or CO<sub>2</sub>-bearing brines also have some quality or consistency problems. Some of the immiscibility boundaries determined by Gehrig coworkers (Gehrig, 1980; Gehrig et al., 1986) seem to have large deviations from other measurements or model predictions. The volumetric data of Gehrig (1980), Nighswander et al. (1989), Teng and Yamasaki (1998), and probably the CO<sub>2</sub> brine densities of Li et al. (2004) as well, are believed to have systematic deviations. Therefore, we still lack the sufficient and accurate volumetric data for the ternary to adequately model CO<sub>2</sub> sequestration.

## 5.2. Thermodynamic modeling of CO<sub>2</sub>–H<sub>2</sub>O and CO<sub>2</sub>–H<sub>2</sub>O–NaCl fluids

(i) Most models are developed for homogeneous fluids; very few are for phase equilibria. As mentioned in the introduction, an accuracy better than 0.1% is necessary for a density model in  $CO_2$  sequestration modeling. However, none of the existing models is sufficient for this purpose, although there are a few models capable of predicting the liquid densities with accuracies close to those of experiments, such as the models of Duan et al. (1992b), Ji et al. (2005), Song et al. (2003b), O'Connell et al. (1996), Sedlbauer et al. (2000), and Plyasunov et al. (2000a,b).

Currently, establishing an accurate binary density model over a wide *PTx* range should be possible if the model is well constructed and its parameters are constrained with the best data sets in existence, such as those of Ohsumi et al. (1992), Song et al. (2005), Crovetto and Wood (1992), Hnedkovsky et al. (1996), Seitz and Blencoe (1997), Singh et al. (2000), Blencoe et al. (2001a), Li et al. (2004), and Hebach et al. (2004) (if the systematic deviations are corrected). As for the ternary, it is possible to develop a predictive density model using the existing ternary data plus the associated subsystem data if the model is based on a sound theory.

- (ii) The existing equations of state are still inadequate for the PVTx prediction in and near critical regions. This failure results from the nonclassical (singular) behavior of fluids near their critical points, such as the diversity of  $[\partial V/\partial P]_{T,x}$ . This problem cannot be rigorously solved in the framework of the classical (analytical) EOS based on mean field theory, such as the numerous cubic, virial, SAFT and perturbation theory EOS, etc. A strict solution to this problem requires a nonclassical EOS, which must contain a non-analytical term to account for the singularity in the near-critical region. Now, many approaches have been developed for the accurate representation of near-critical behavior of fluids (Wilson, 1971; Chang and Doiron, 1983a; Wilson, 1983; Woolley, 1983; Sengers and Levelt Sengers, 1984; Anisimov et al., 1985; Albright et al., 1986; Sengers, 1986; Rainwater, 1989; White and Zhang, 1993; Anisimov et al., 1995; Kiselev, 1997; Edison et al., 1998; Lue and Prausnitz, 1998; Tang, 1998; Kiselev and Ely, 1999; Kiselev and Friend, 1999; White, 2000; van't Hof et al., 2001; Hu et al., 2003a,b; Kiselev and Ely, 2004). Almost all these approaches are based on the so-called renormalization group theory and crossover functions, which involve complicated knowledge of mathematics. Different from these approaches, the fluctuation solution theory is a very useful tool for us to solve the near-critical singularity problem. However, this theory is only valid for dilute solutions. Now, it is still an unsolved problem how to develop an equation that has a close form but can accurately describe both classical and nonclassical behaviors of fluids.
- (iii) For the ternary system, only a few simple density models were proposed to correlate experimental data at low temperatures below 333.15 K, with accuracy close to that of experiments. At higher temperatures, there is no accurate model for the PVTx properties. Because of the great differences in compositions and physicochemical properties, accurate PVTx prediction of coexistent phases with a single model is very difficult. For the liquid PVTx modeling, the difficulty comes from the hydrogen bond, hydration, reaction, ionization and association phenomena of NaCl, CO<sub>2</sub>, and H<sub>2</sub>O. In this situation, a hybrid model may be an effective option. For example, one can use an equation of state for the gaseous phase, and an activity coefficient model for liquid phases, or use independent parameters for each of the coexistent phases. Using excess volume or apparent molar volume models is also a good option, because this approach has no constraint on the relevant subsystem models.

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